

Magnetism

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Optical Condensed Matter Physics

Lectures

Last time	Introduction	Ch.1; 2.1-2.5; 8.9
Today	Moments, interactions, environment	3.1, Ch.4; 7.1-7.7

Magnetism

Origin of magnetism

- Orbital magnetism ($p \rightarrow p+eA$)
- Spin magnetism (spin S , localized, itinerant)

Magnetic interactions

- Dipole
- Exchange interaction (direct, indirect)
- RKKY (itinerant exchange)
- Double exchange
- Anisotropic exchange (LS interaction)

Magnetism

Diamagnetism:

- No magnetic moments
- No magnetic interaction
- Response due to induced currents
- Magnetization opposite to field

Paramagnetism:

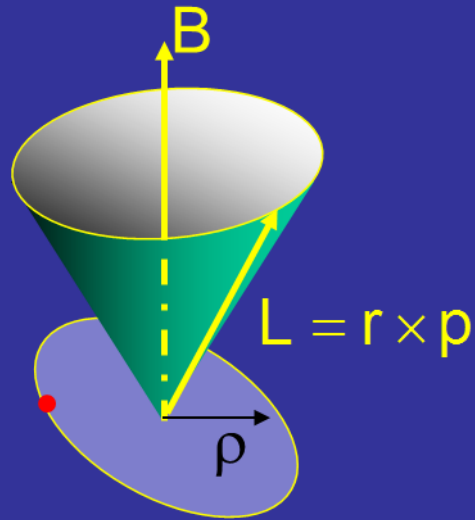
- Magnetic moments (spin, orbit)
- Weak magnetic interactions
- Response due to orientation
- Magnetization in field direction

Ordered magnetism:

- Magnetic moments
- Strong magnetic interactions
- Response due to polarization
- Ferro-, antiferro-, ferrimagnetic

'Non-ordering magnetism'

'Classical' Langevin diamagnetism



e^- in B-field \Rightarrow Larmor precession with

$$\omega_L = \frac{eB}{2mc}$$

Current of Z electrons per atom

$$I = -Ze \cdot \frac{\omega_L}{2\pi}$$

Magnetic moment

$$\mu = I \cdot A = I \cdot \pi \langle \rho^2 \rangle = -\frac{Ze^2 B}{4mc^2} \cdot \langle \rho^2 \rangle$$

Susceptibility of n atoms/volume

$$\chi = \frac{n\mu}{B} = -\frac{nZe^2}{6mc^2} \cdot \langle r^2 \rangle$$

Force:
$$\mathbf{F} = M \frac{\partial H}{\partial \mathbf{x}}$$

Floating matter (magnetic levitation)

(perfect) diamagnetism

Floating YBaCuO_{6+d}



1986: J.G. Bednorz
& K.A. Mueller



Magnetic moments

- Intrinsic magnetic moment of elementary particles: **Spin**
- Moving charges: **Orbits**

- Spin-orbit coupling $\rightarrow J; m_J$
- Magnetic moment:

$$\mu = g_J \sqrt{J \cdot (J + 1)} \mu_B$$

$$\mu_z = g_J \cdot m_J \cdot \mu_B$$

- Energy in a field

$$E = g_J \cdot m_J \cdot \mu_B \cdot B$$

- Free electron spin ($S=1/2$) in 1 T:

$$\mu_s \approx 1.73 \mu_B \quad E \approx 60 \mu\text{eV} \therefore 0.7 \text{ K}$$

Only orbits: $g_l=1$

Only electron spin: $g_s=2$

More general:

$$g_J = \frac{3J(J+1) + S(S+1) - L(L+1)}{2J(J+1)}$$

Bohr magneton

$$\begin{aligned} \mu_B &= \frac{\hbar e}{2m_e} = 9.274 \text{ J/T} \\ &= 5.7553 \cdot 10^{-5} \text{ eV/T} \end{aligned}$$

Localized moments: partially filled shells

- d-levels, f-levels
- L-S coupling scheme
- Hund's rules
 1. **Maximize S** (\rightarrow spatially antisymmetric, decrease shielding)
 2. **Maximize L** (\rightarrow orbiting in same direction, decrease repulsion)
 3. **$J=|L-S|$ for $<$ half full**
 $J=|L+S|$ for $>$ half full (\rightarrow minimize LS interaction)

example Fe^{2+} : $[\text{Ar}]3d^6$

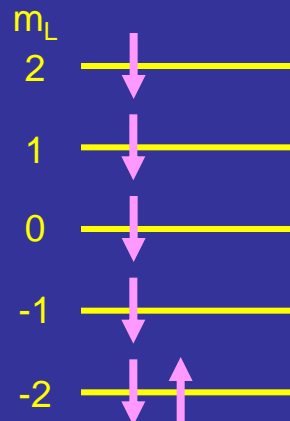
6 electrons

Total spin: 2

Total L: 2

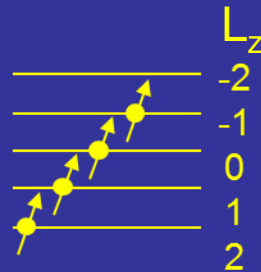
$\rightarrow J=4$

Term symbol ($(^{2s+1})L_J$): 5D_4



Some more examples

d-shell ($L=2$), 4 electrons
 $\text{Mn}^{3+}, \text{Cr}^{2+}$



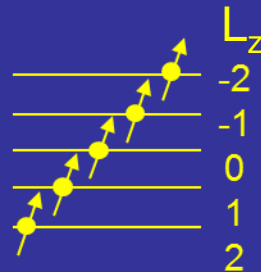
$$S=2$$

$$L=2$$

$$J=0$$

$$\Rightarrow {}^5D_0$$

d-shell ($L=2$), 5 electrons
 $\text{Fe}^{3+}, \text{Mn}^{2+}$



$$S=5/2$$

$$L=0$$

$$J=5/2$$

$$\Rightarrow {}^6S_{5/2}$$

d-shell ($L=2$), 6 electrons
 Fe^{2+}



$$S=2$$

$$L=2$$

$$J=4$$

$$\Rightarrow {}^5D_4$$

Environment: crystal field

Rare earth's: 4f shell's small ('inner' electrons)

Iron group: 3d shell's on the outside

=> decoupling of L and S, J no longer good Quantum number

=> splitting of the $2L+1$ orbital states

=> Quenching of the orbital angular momentum ($L_z \rightarrow 0$)

=> High spin – Low spin transitions

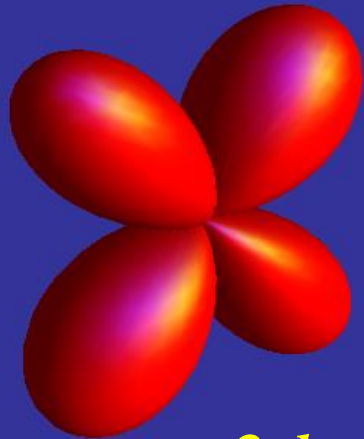
=> Jahn-Teller distortions

=> Orbital excitations (orbitons)

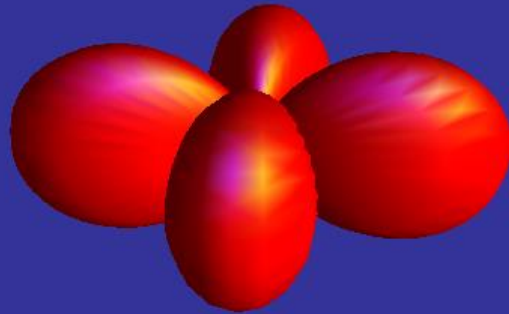
Anisotropy

Kramers degeneracy (local B-field probe)

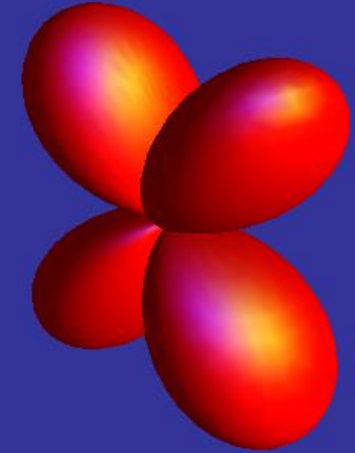
3d orbitals (Chem)



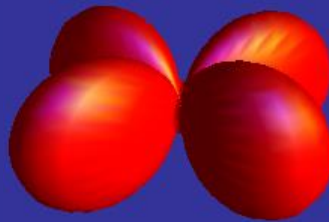
$3d_{xz}$



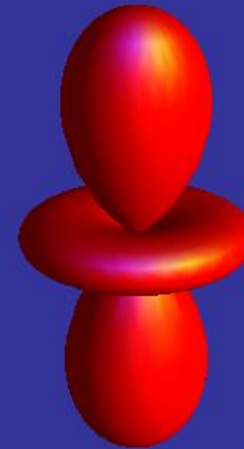
$3d_{xy}$



$3d_{yz}$

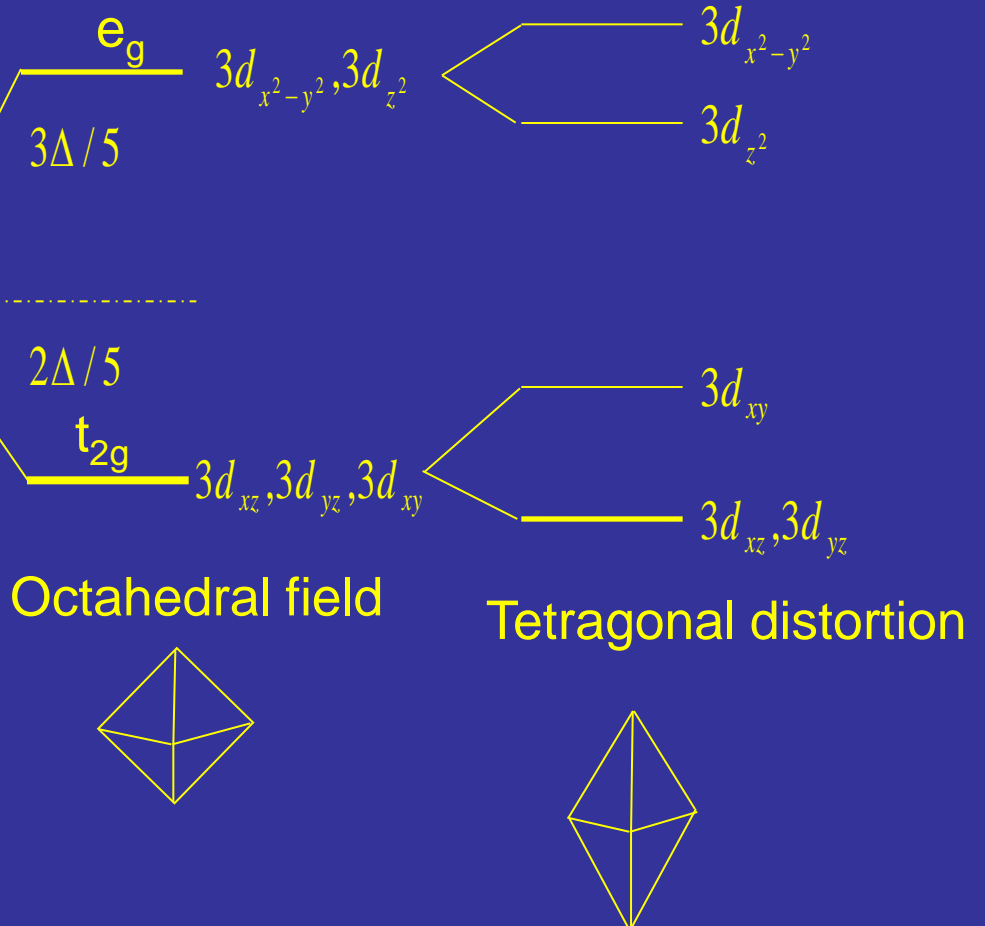
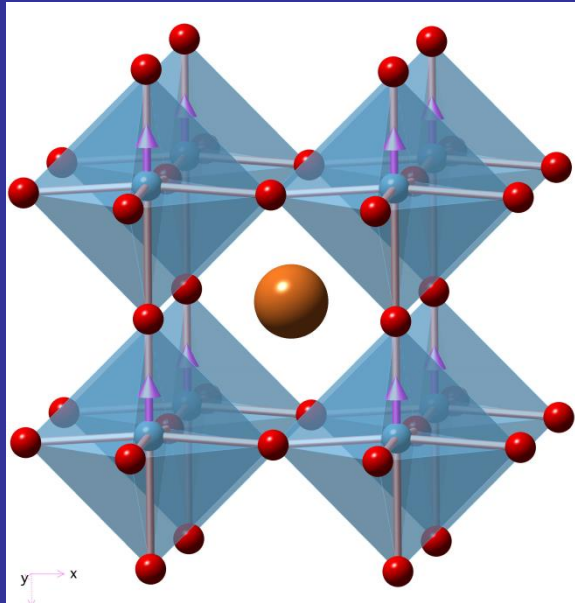
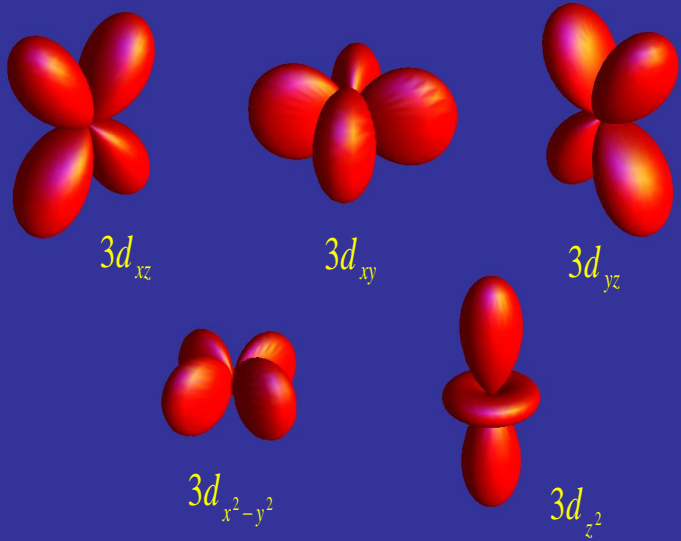


$3d_{x^2-y^2}$

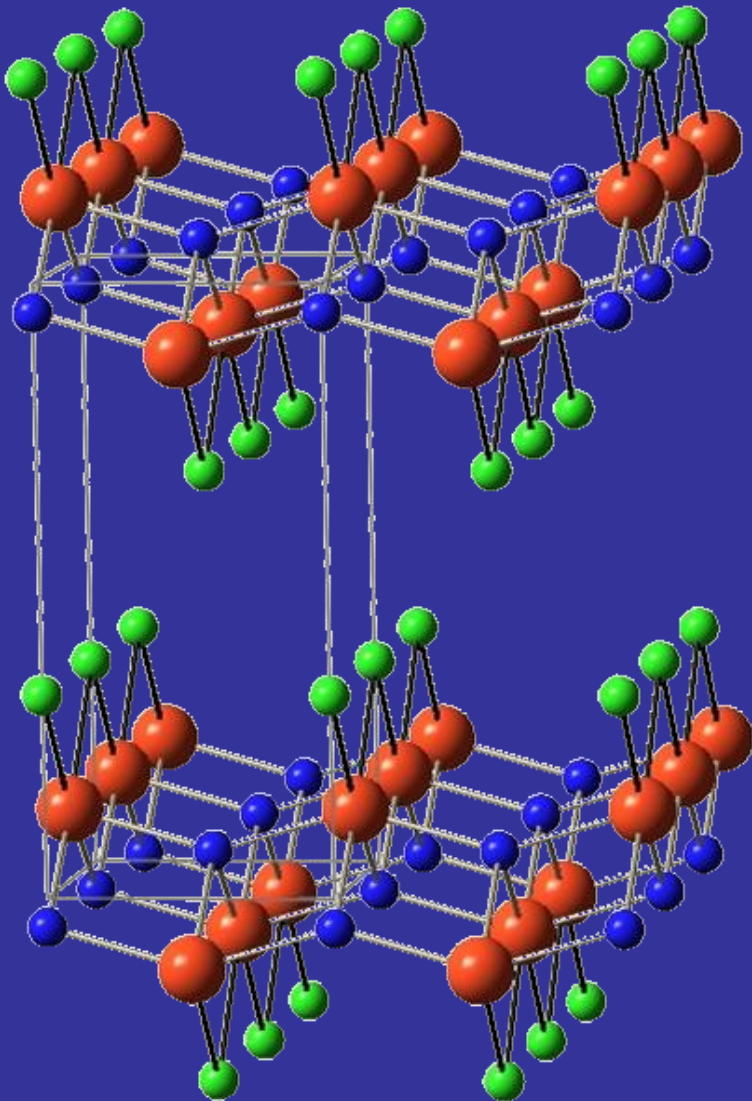


$3d_{z^2}$

3d orbitals in a CF



Perovskite structure (ABO_3 e.g. $BaTiO_3$)



Pmmn

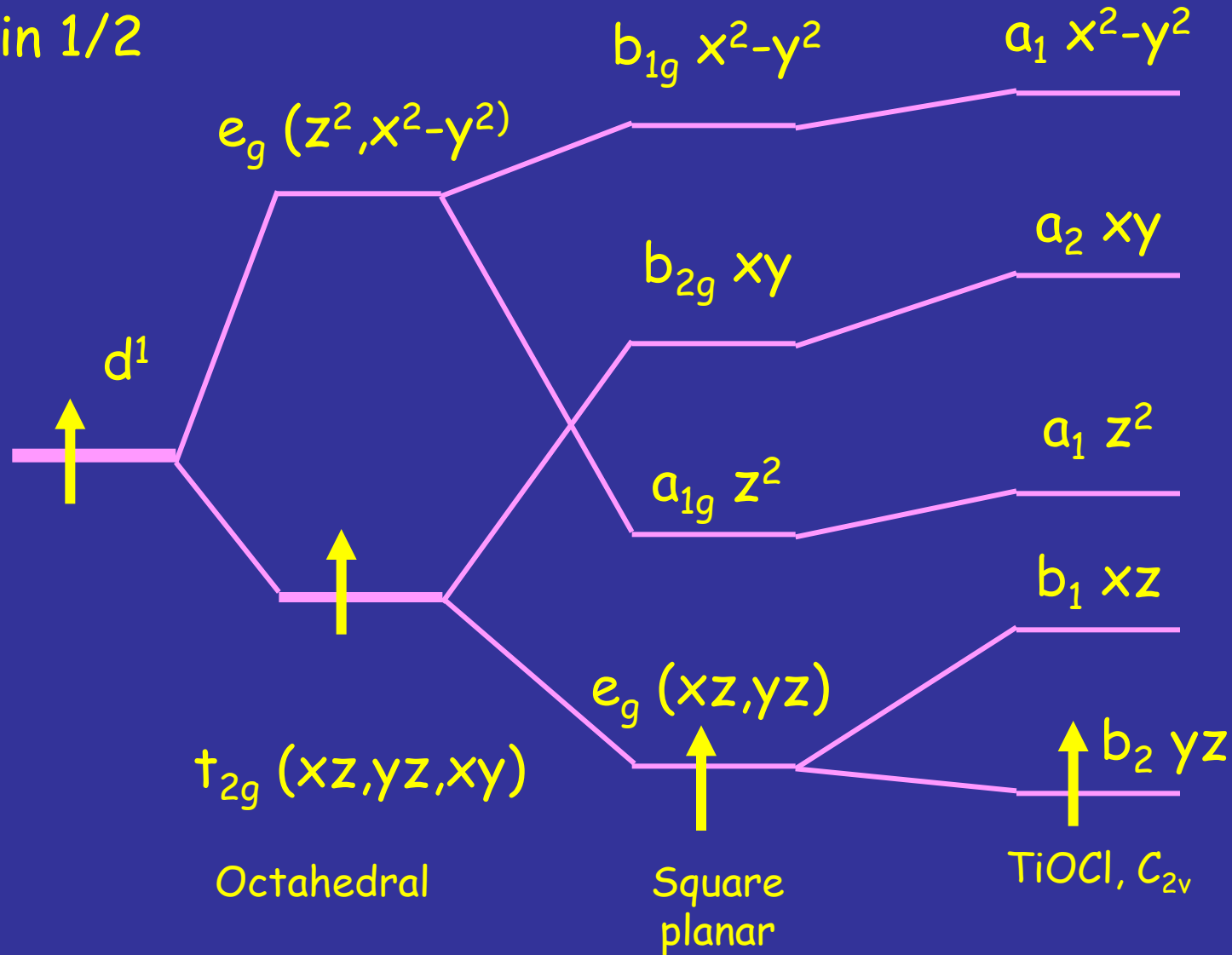
Double bi-layer of
 TiO_4Cl_2 "Octaedra"

Layered:
Cl-Cl distance 3.7 Å

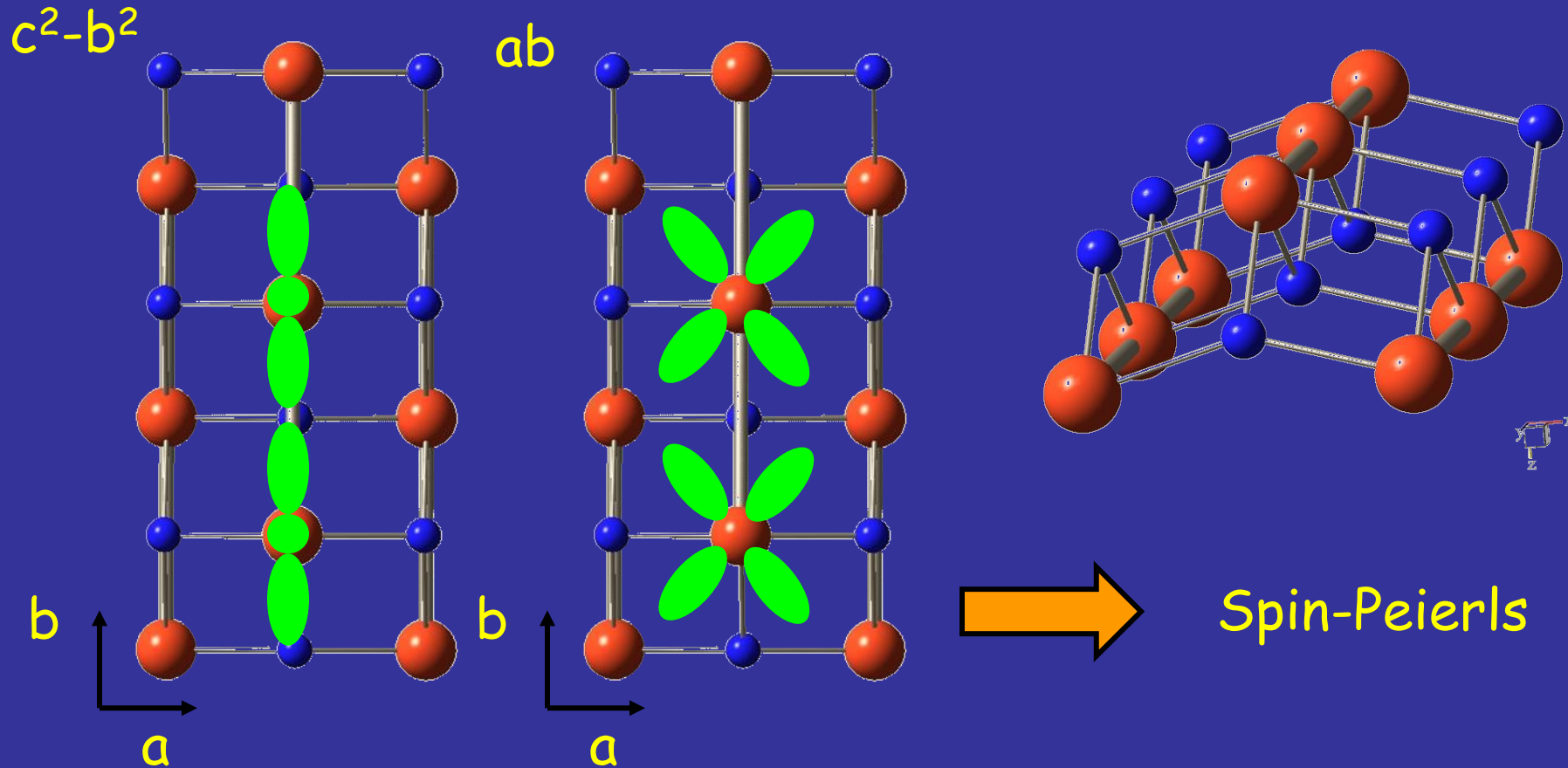


3d levels in TiOX

Ti⁺ 3d¹, spin 1/2

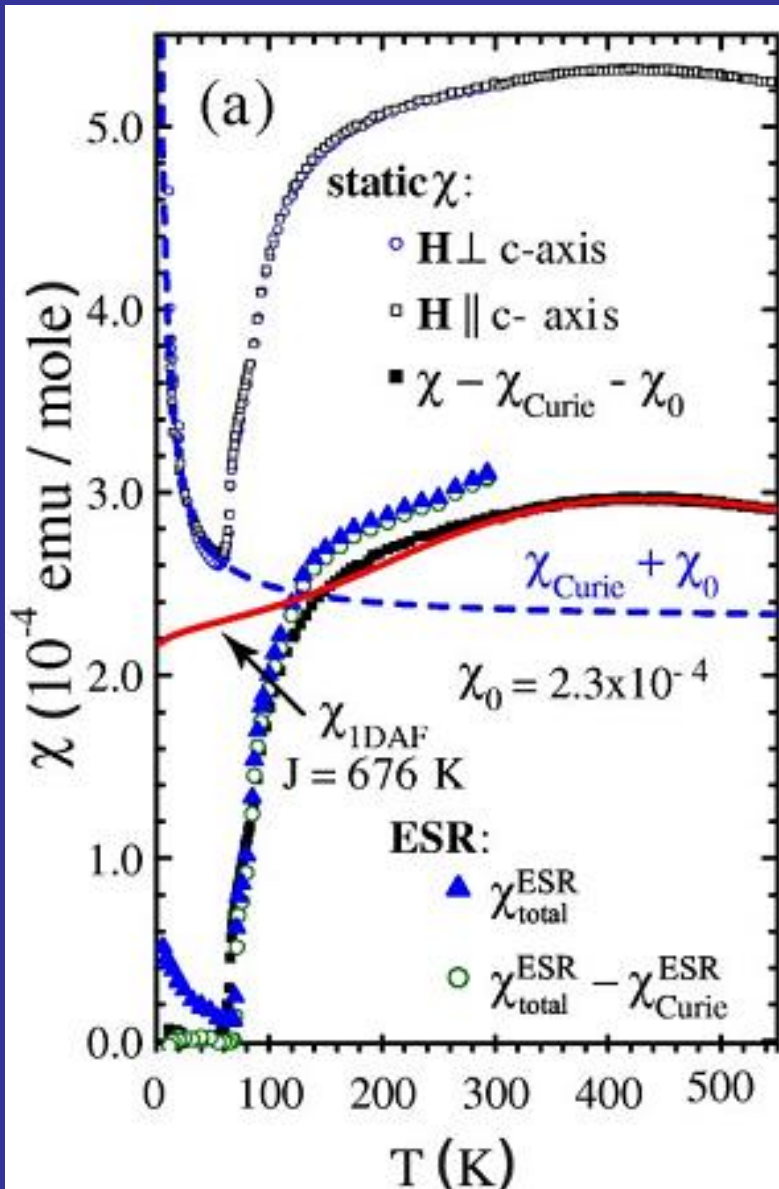


Occupied orbitals



Exchange through direct Ti-Ti interaction along $b \rightarrow$ 1D system
Mott-Hubbard \rightarrow AF interaction

Susceptibility $TiOX$



Above T_c : Typical 1D Heisenberg AF
Below T_c : Opening of a gap due to dimerization

Kataev et al., PRB 68, 140405 (2003)