

Magnetic moments
Types of magnetism

Environment

Rare earth's: 4f shell's small ('inner' electrons)

Iron group: 3d shell's on the outside

=> decoupling of L and S, J no longer good Quantum number

=> splitting of the $2L+1$ orbital states

=> Quenching of the orbital angular momentum ($L_z \rightarrow 0$)

=> High spin – Low spin transitions

=> Jahn-Teller distortions

=> Orbital excitations (orbitons)

Anisotropy

Lectures

Lect. 1	Introduction	Ch.1; 2.1-2.5; 8.9
Lect. 2	Interactions, environment	3.1, Ch.4; 7.1-7.7
Lect. 3	Ordering, Domains	5.1-5.3; 6.7; 8.3, 8.7, 8.8
Lect. 4	Symmetry breaking	6.1-6.6
Lect. 5	Quantum magnetism	8.1-8.6

- Dipole – Dipole
- Direct exchange (H_2 molecule)
- Indirect exchange
- Double exchange
- Anisotropic exchange
- Rudeman Kittel Kasuya Yoshida (RKKY)
- Stoner (“spontaneous Pauli”)

- Ch. 4 & 7.2,7.3,7.7

Dipole-dipole interaction

Dipole-dipole interaction is an anisotropic interaction

$$E = \frac{\mu_0}{4\pi} \frac{\vec{\mu}_1 \cdot \vec{\mu}_2 - 3(\vec{\mu}_1 \cdot \vec{e}_{12})(\vec{\mu}_2 \cdot \vec{e}_{12})}{r_{12}^3}$$

$$E \approx 10^{-23} \text{ J} \sim 1 \text{ K} \text{ for } r = 2 \text{ \AA} \text{ and } \mu = \mu_B$$

In real materials: $T_c \sim 10^2 - 10^3 \text{ K} !!$

→ Dipole-dipole interaction hardly ever dominates

H₂ molecule

- LCAO gives wrong solution (e.g. triplet ground state)
- Correlated picture (Heitler-London approach) is better
 - Starting point: two electron orbitals $|\phi_a(1)\phi_b(2)\rangle$ and $|\phi_a(2)\phi_b(1)\rangle$
 - Wavefunctions

$$\Psi_s = \frac{1}{\sqrt{2}} [\phi_a(1)\phi_b(2) + \phi_a(2)\phi_b(1)] \chi_s$$

$$\Psi_t = \frac{1}{\sqrt{2}} [\phi_a(1)\phi_b(2) - \phi_a(2)\phi_b(1)] \chi_t$$

- Spin parts:

$$\chi_s = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

$$\chi_t = \begin{cases} |\uparrow\uparrow\rangle \\ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \\ |\downarrow\downarrow\rangle \end{cases}$$

H₂ molecule

- Singlet ground state
- Exchange energy: Pauli + Coulomb

$$H = -\frac{\hbar^2}{2m} \left(\nabla_1^2 + \nabla_2^2 \right) + V_c(1,2)$$

$$V_c(1,2) = \frac{e^2}{|r_1 - r_2|} + \frac{e^2}{|R_1 - R_2|} - \frac{e^2}{|r_1 - R_2|} - \frac{e^2}{|r_2 - R_2|}$$

$$E_S - E_T = 2 \langle \phi_a(1) \phi_b(2) | V_c | \phi_a(2) \phi_b(1) \rangle := 2J$$

$$H = -2J \vec{S}_1 \cdot \vec{S}_2 = \begin{cases} -2J \cdot -\frac{3}{4} & \text{for the spin singlet} \\ -2J \cdot \frac{1}{4} & \text{for the spin triplet} \end{cases}$$

Exchange interaction

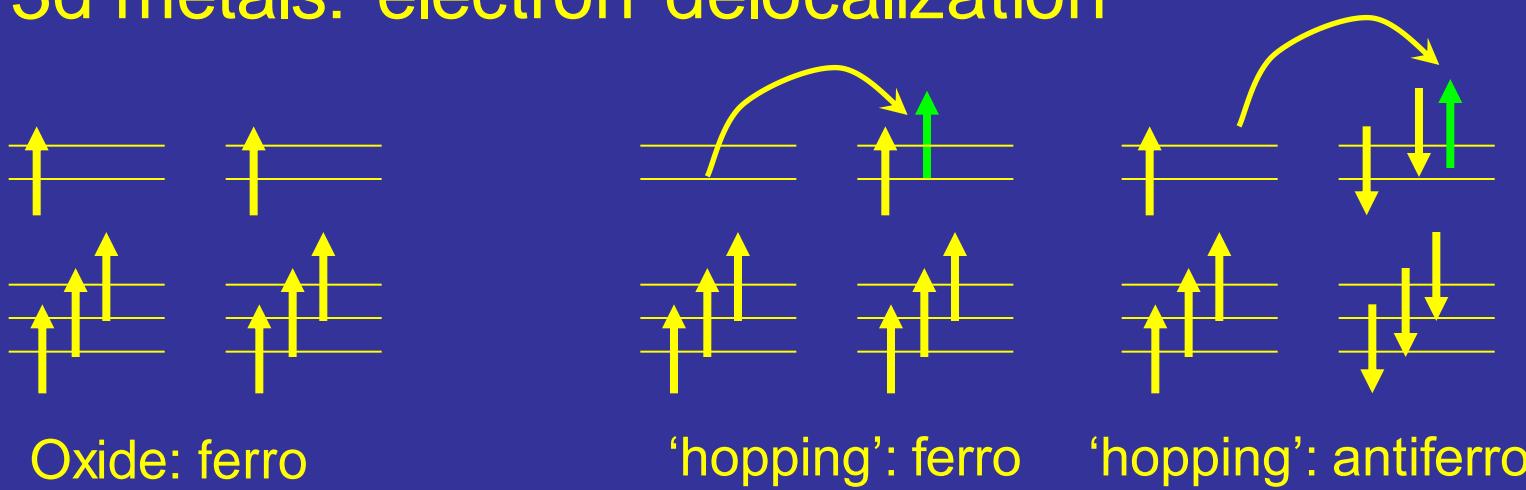
- Heisenberg Hamiltonian

$$H = -2 \sum_{i,j} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

- $J>0$: Ferro
- $J<0$: Antiferro

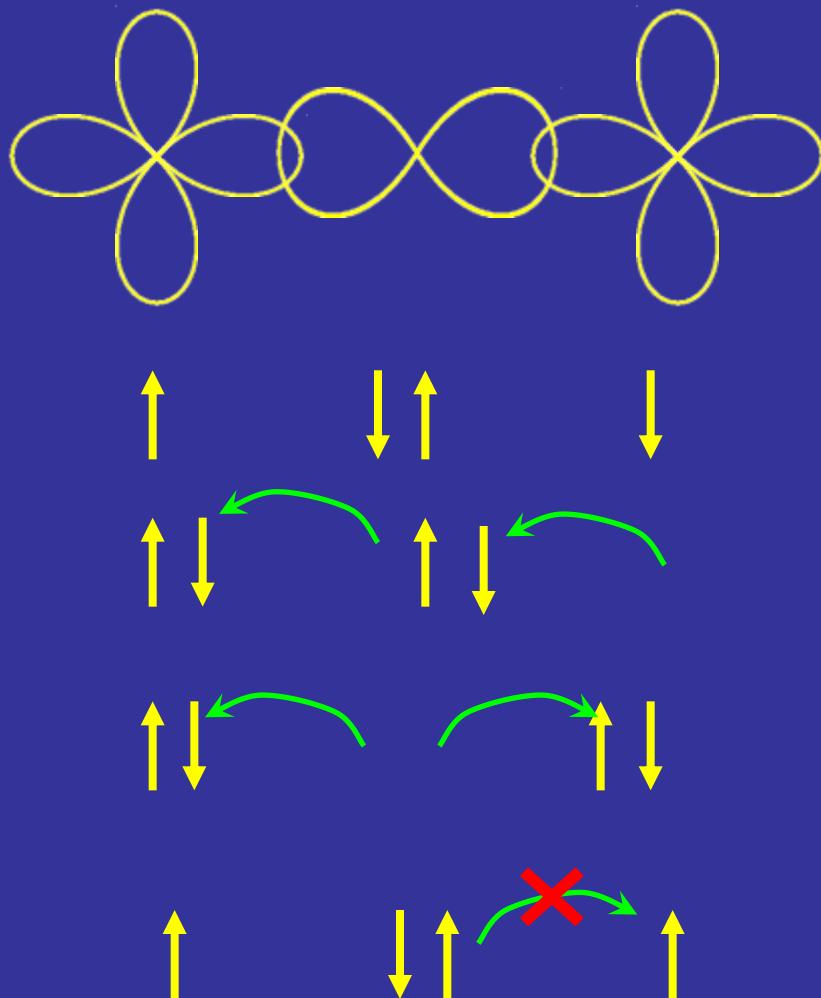
Direct exchange

- Need direct wave function overlap
 - Often ferromagnetic
 - Small in 4f, 5f elements
 - Can be important in 3d oxides (but see indirect!)
 - In 3d metals: electron delocalization



- Relatively small (but remember TiOX)
- Depends on orbital occupation and geometry

Indirect exchange



Ground state antiferro

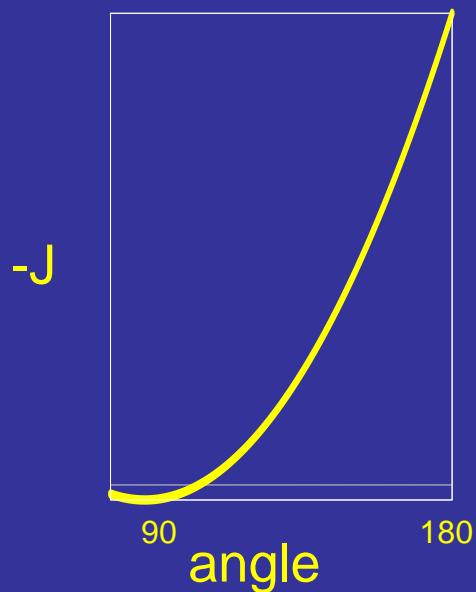
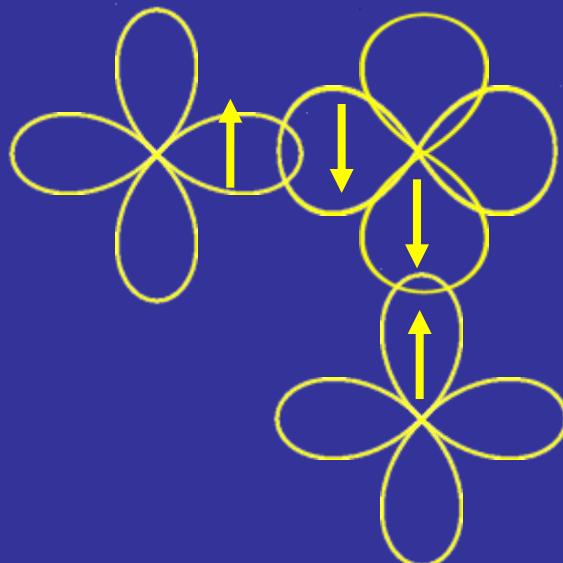
'2*hopping'

Pauli forbidden

Hopping → delocalization → energy gain
Energy: 2 hops = $2t$; cost = U
 $\rightarrow J \sim -t^2/U$

Examples: High T_c's; MnO; MnF₂

Indirect exchange



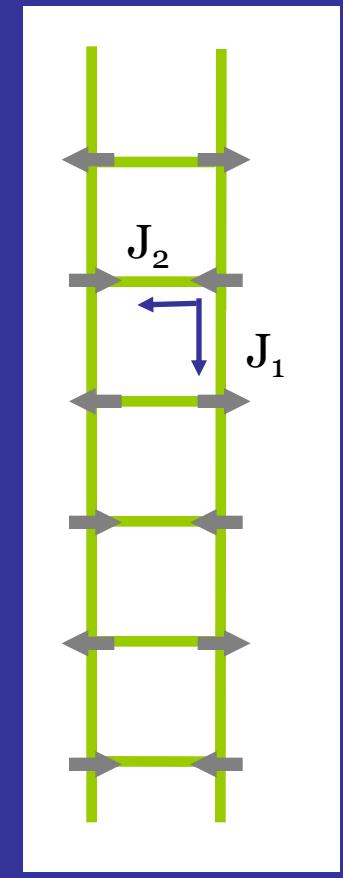
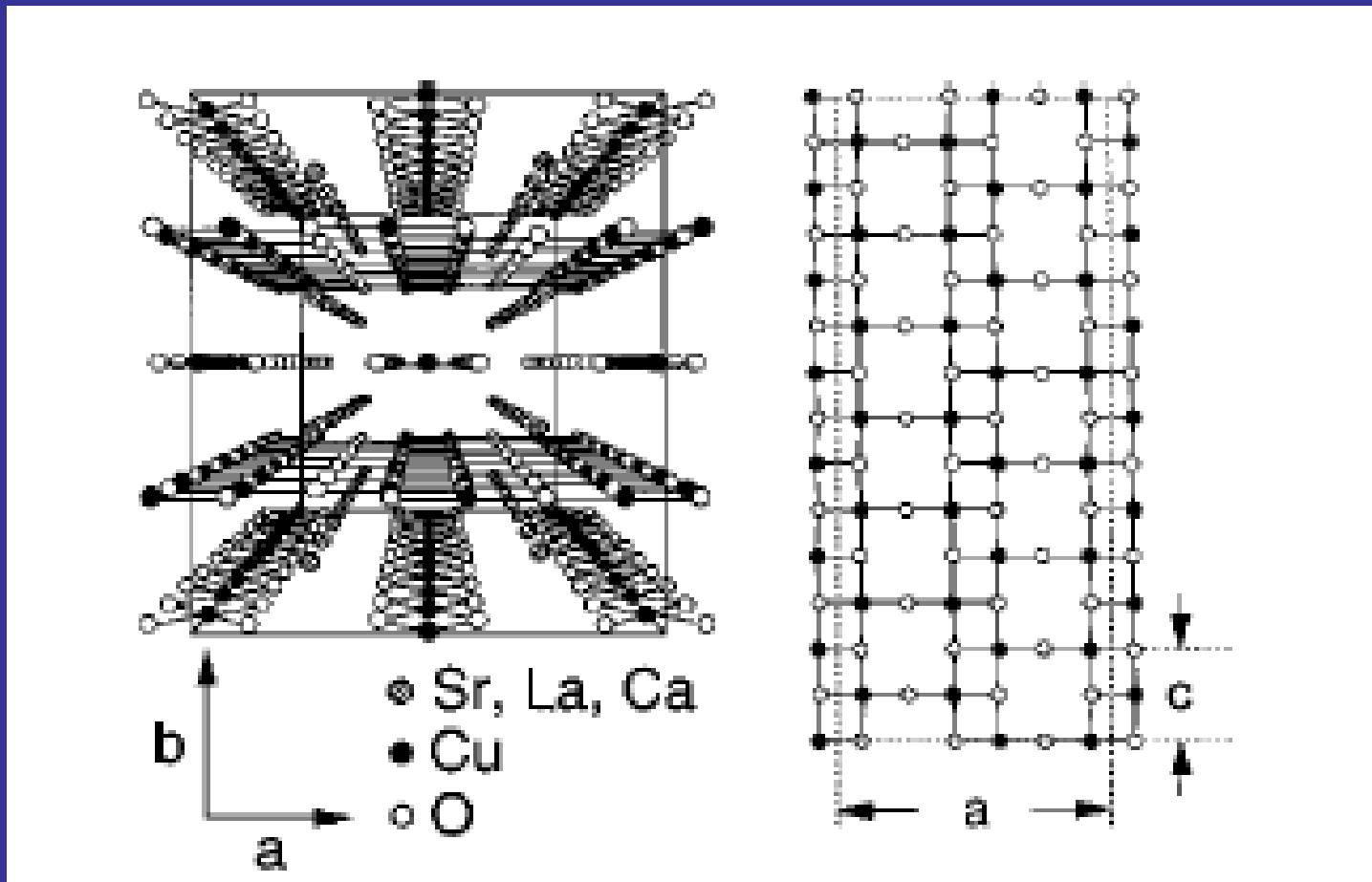
Hopping \rightarrow delocalization \rightarrow energy gain
Energy: 2 hops = $2t$; cost = U

$$\Rightarrow J \sim -t^2/U$$

Examples: High T_c 's; MnO ; MnF_2 ;
telephone number compound

Relatively strong (depends on U)
Usually AF (F when not same 3d, e.g. d^3-d^5)
Strongly dependent on angle of bonding
at 180° strongly AF
at zero weakly F
(Goodenough-Kanamori rules)

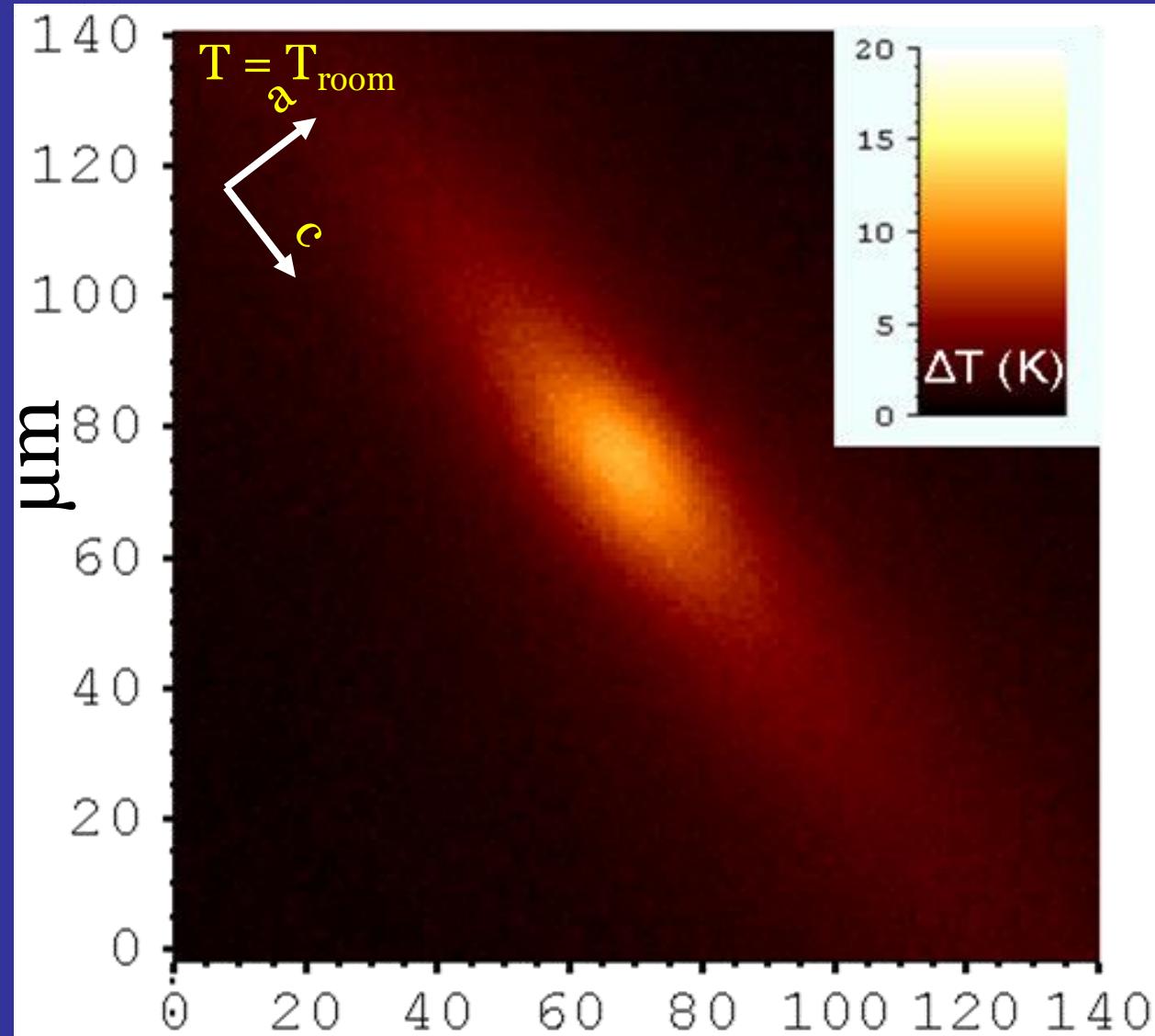
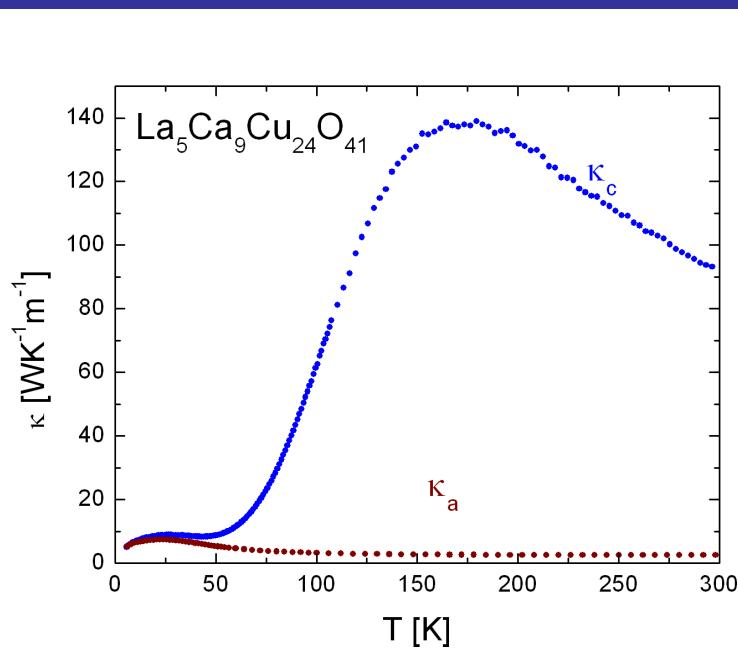
$(Sr,La,Ca)_{14}Cu_{24}O_{41}$



$J_1 = 130 \text{ meV}$
 $J_2 = 70 \text{ meV}$
 $\Delta = 32 \text{ meV}$

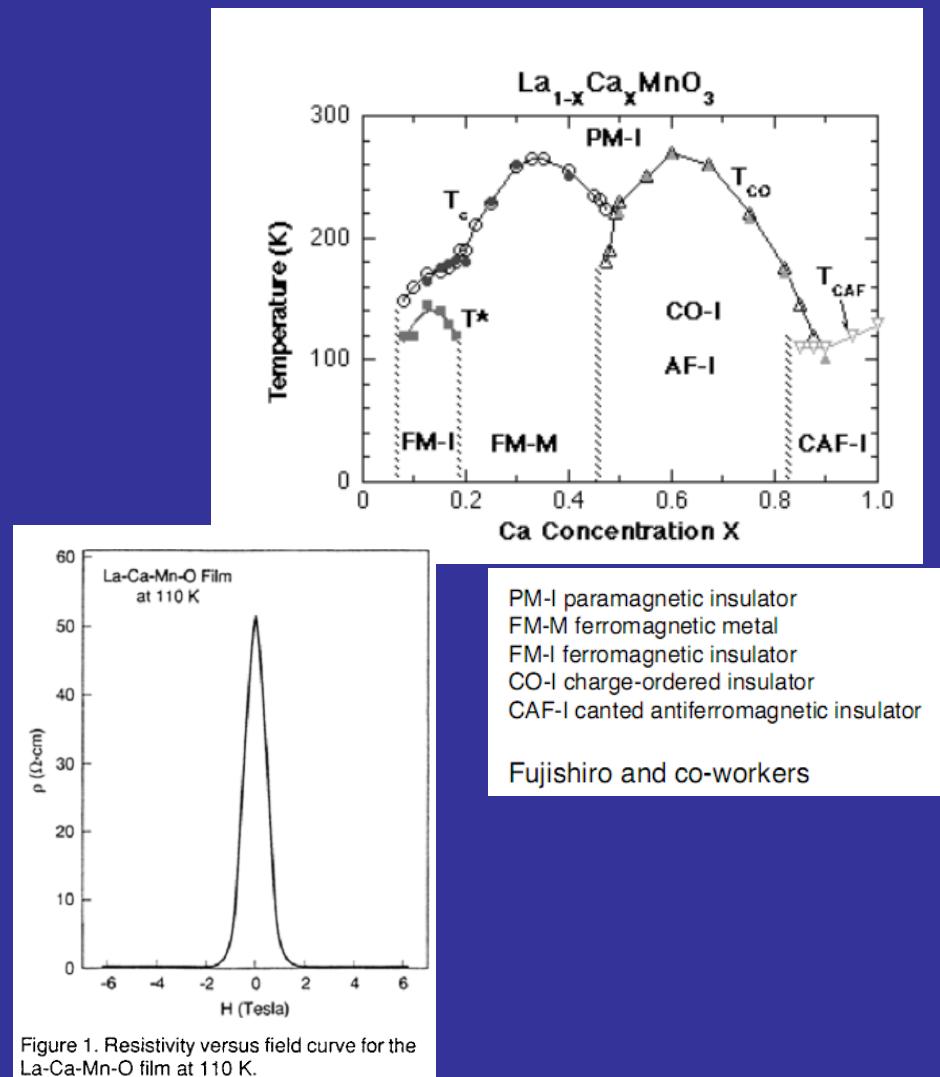
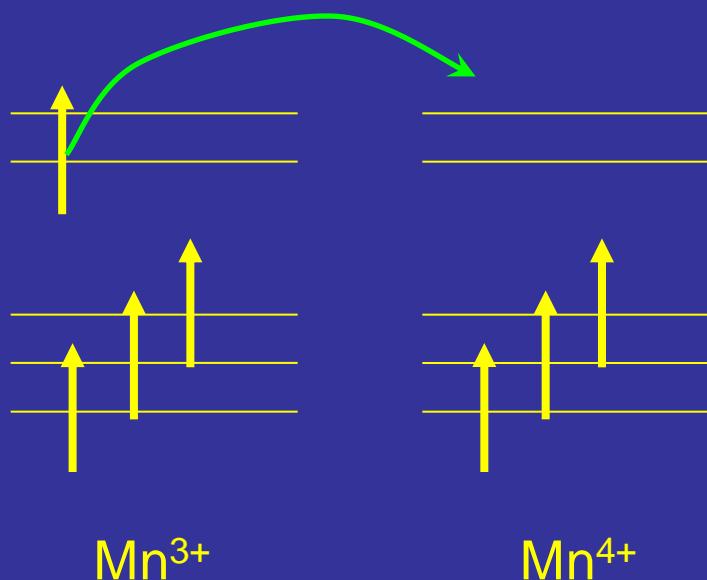
Eccleston *et al.*, PRL 81, 1702 (1998)

$La_9Ca_5Cu_{24}O_{41}$



Double exchange

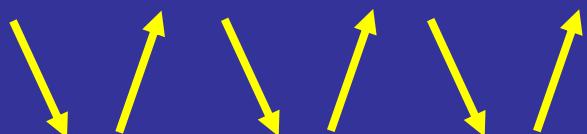
- Mixed valence
- Usually ferro metal
- Relatively strong
- $\text{La}_{1-x}\text{Sr}_x\text{MnO}_3$ ($\text{Mn}^{3+}/\text{Mn}^{4+}$)
 Fe_3O_4 (AB_2O_4 , $\text{Fe}^{2+}/\text{Fe}^{3+}$)



CMR in para phase close to T_{curie}

Anisotropic exchange

- Dzyaloshinsky-Moriya interaction
- Mixing in of excited d-states
- LS coupling in excited state
- If spins inversion symm. related then 0
- Form different from Heisenberg: $\vec{D} \cdot \vec{S}_i \times \vec{S}_j$
- Favors perpendicular alignment
- Examples: $\alpha\text{-Fe}_2\text{O}_3$, MnCoO_3
- In AF's leads to net moment \rightarrow weak ferro (canted AF)



Effects of Dzyaloshinskii-Moriya interaction

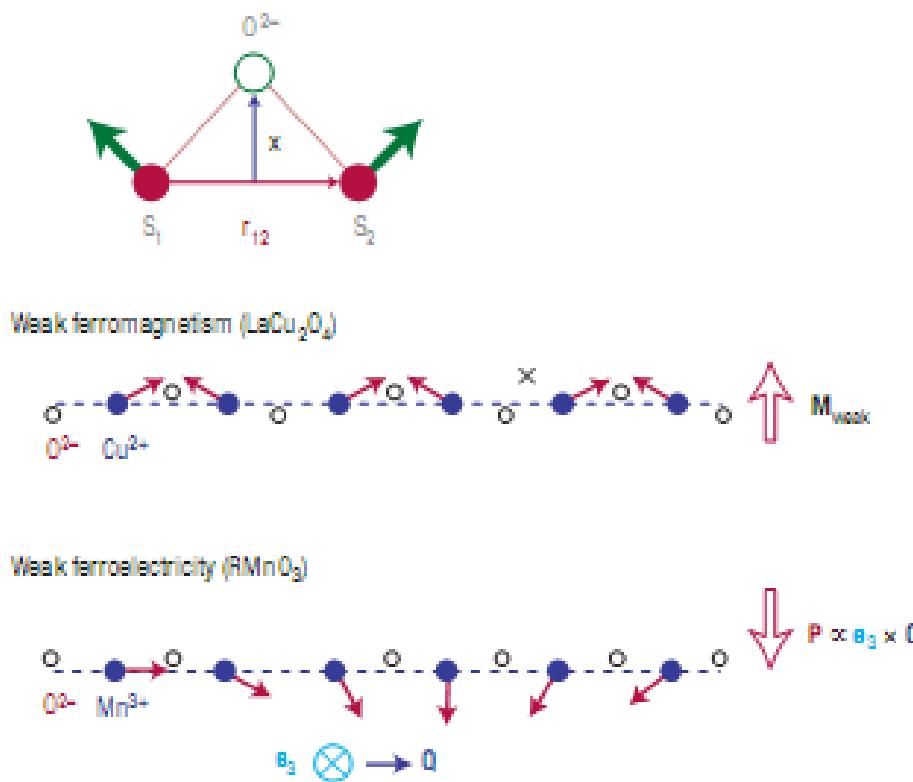
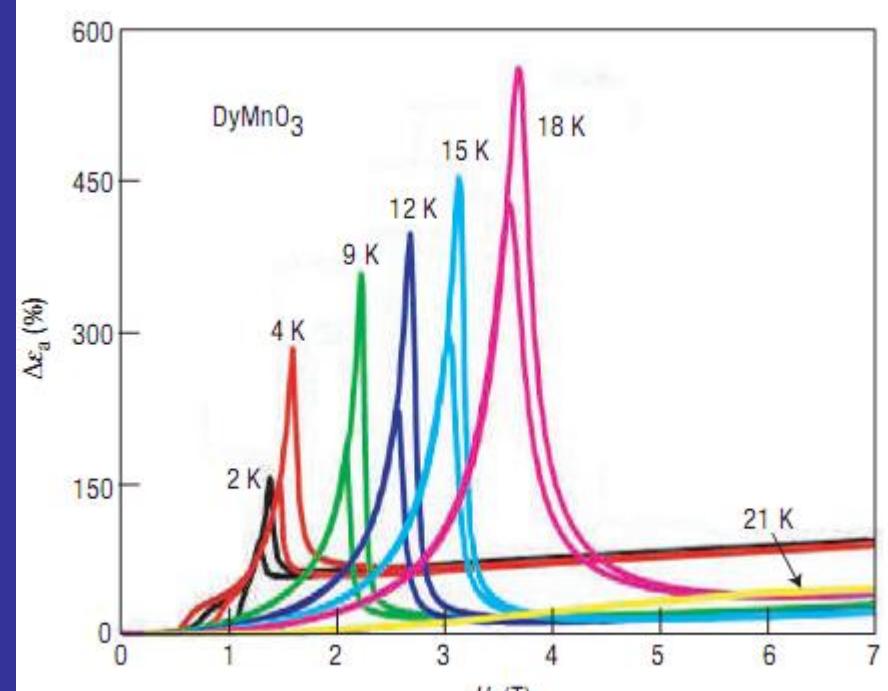


Figure 5 Effects of the antisymmetric Dzyaloshinskii–Moriya interaction. The interaction $H_{DM} = D_{12} \cdot [S_1 \times S_2]$. The Dzyaloshinskii vector D_{12} is proportional to spin-orbit coupling constant λ , and depends on the position of the oxygen ion (open circle) between two magnetic transition metal ions (filled circles), $D_{12} \propto \lambda x \times \hat{r}_{12}$. Weak ferromagnetism in antiferromagnets (for example, LaCu_2O_4 layers) results from the alternating Dzyaloshinskii vector, whereas (weak) ferroelectricity can be induced by the exchange striction in a magnetic spiral state, which pushes negative oxygen ions in one direction transverse to the spin chain formed by positive transition metal ions.



$$\mathbf{P} \propto [(\mathbf{M} \cdot \partial) \mathbf{M} - \mathbf{M} (\partial \cdot \mathbf{M})].$$

Mostovoy & Cheong

Excercises

- 3.1; 3.5 ; 3.11
- Which 3d ions can be Jahn Teller active, what are the corresponding d-electron occupations ?
- The lowest 3d levels in an octahedral field are the t_{2g} triplet (states like d_{xy}). Can you predict the energy splitting in a tetrahedral environment
- 4.4; 4.5; 5.1; 6.13