

Condensed Matter Physics I

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Previously

Thermal expansion

- Lattice anharmonicity

Heat Capacity

- Dulong-Petit law
- Density of states
- Debye model (T^3 law)
- Einstein model
- Phase transitions

Today

- Heat conductivity
- Intro to second quantization
- Metals

Thermal conductivity



Thermal energy flux: $J_u = \frac{dw}{dt}$

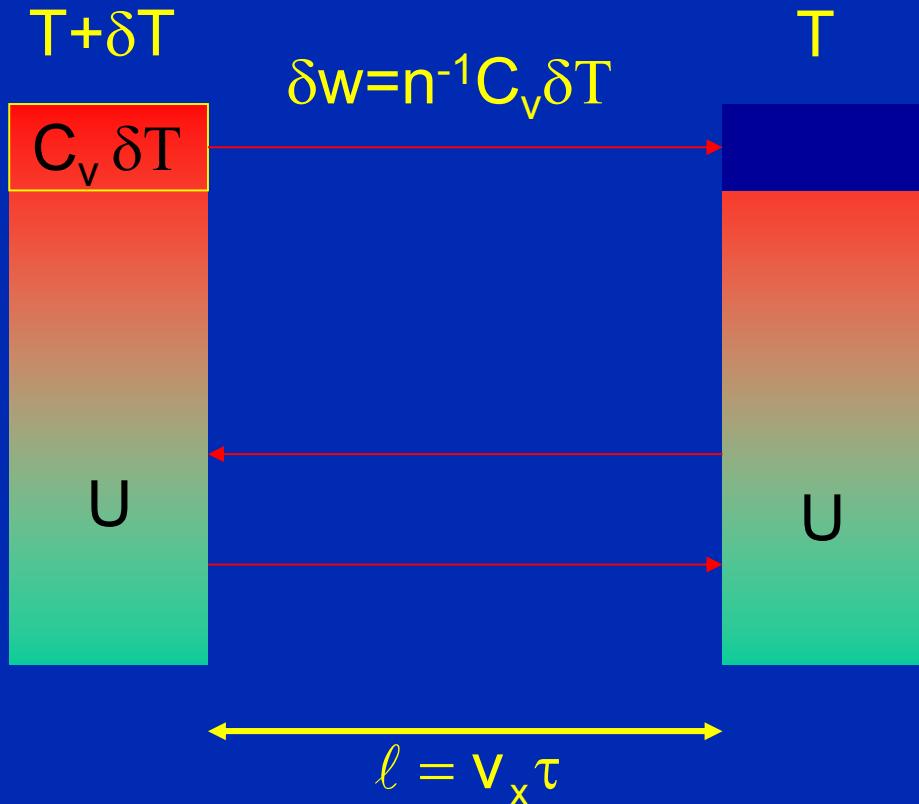
Thermal gradient: $\frac{dT}{dx} = \frac{T_r - T_l}{L}$

$$J_u = -\kappa \frac{dT}{dx}$$

κ : Thermal conductivity
Is due to diffusion of particles (here phonons)

Thermal conductivity

$$J_u = -\kappa \frac{dT}{dx}$$



$$J_u = \frac{dw}{dt} = n \langle \mathbf{v}_x \cdot \delta \mathbf{w} \rangle = C_v \langle \mathbf{v}_x \delta T \rangle$$

$$\delta T = -v_x \tau \cdot \frac{dT}{dx}$$

$$J_u = -C_v \tau \frac{dT}{dx} \langle v_x^2 \rangle$$

$$\left. \begin{aligned} \langle v_x^2 \rangle &= \frac{1}{3} v^2 \\ \tau &= \ell / v \end{aligned} \right\} \Rightarrow J_u = -\frac{1}{3} C_v v \ell \frac{dT}{dx}$$



$$\kappa = \frac{1}{3} C_v v \ell$$

Scattering of phonons

Phonon mean free path ℓ limited by

- Imperfections, crystal boundaries, isotopes
- Phonon-phonon scattering (normal processes)
- Phonon-phonon scattering (Umklapp processes)

Harmonic approximation in perfect lattice:

Phonons are eigenstates of Hamiltonian
NO phonon-phonon scattering

Scattering of phonons

Imperfections etc.: $\ell = T$ independent: $\kappa \propto v\ell C_v(T) \propto T^3$

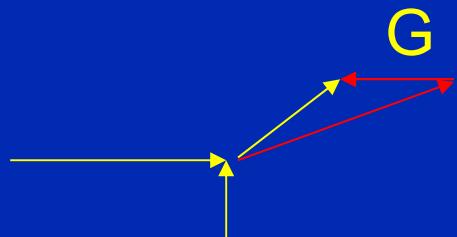
Normal processes: $\hbar\omega_{k_1} + \hbar\omega_{k_2} = \hbar\omega_{k_3}$



No net change of phonon momentum

Umklapp processes: $\hbar\omega_{k_1} + \hbar\omega_{k_2} = \hbar\omega_{k_3}$

$$\vec{k}_1 + \vec{k}_2 = \vec{k}_3 + \vec{G}$$



Momentum transferred to crystal: $\hbar\vec{G}$

Thermal conductivity

$$\kappa = \frac{1}{3} C_v v \ell$$

High T: C_v constant, scattering $\propto n(\omega) \propto T$ $\longrightarrow \kappa \propto T^{-1}$

Low T: Umklapp needs phonon at BZB (\sim Debye energy)

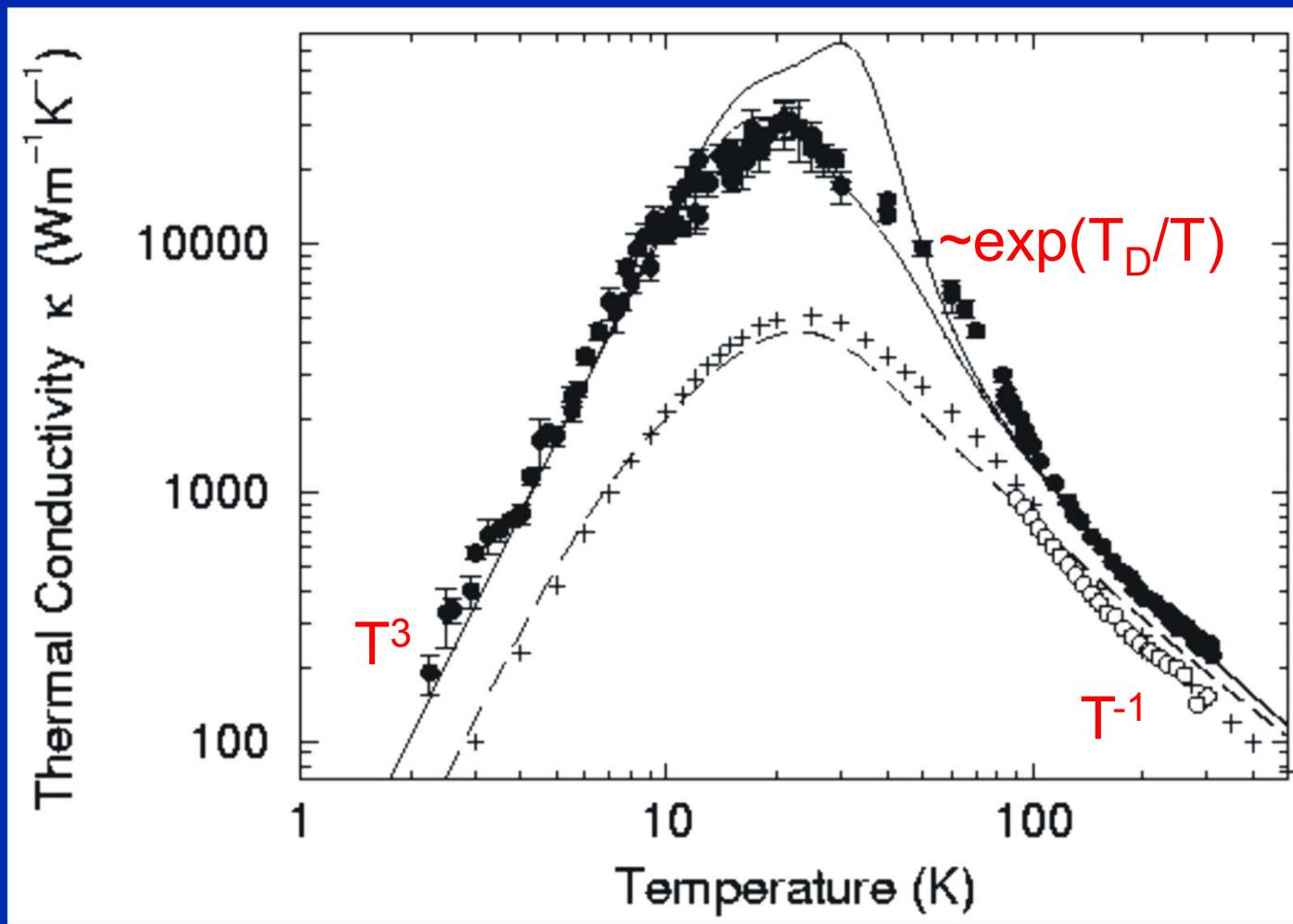
Very low T: no phonons for U-process:
 ℓ constant (imperfections) $\longrightarrow \kappa \propto T^3$

Near $T \sim T_D$ U-processes become important
Scattering $\sim \#$ phonons with \sim Debye energy

$$\ell^{-1} \propto n(\hbar\omega_D/2) = \frac{1}{e^{\theta_D/2T} - 1} \approx e^{-\theta_D/2T}$$

$$\kappa \propto e^{\theta_D/2T}$$

Thermal conductivity



Isotope pure Silicon (T. Ruf et al. Sol. St. Comm. 2000)