## **SUPERCONDUCTIVITY WS 15-16**

Monday 10:00-11:30 SR Exp. physics II

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# Kitchen refrigerator

Expansion Evaporator Coll Compressor Latent Heat  $T=5^{0}C$ of Vaporization COOL GAS Evaporation Compression 2-3 atm COOL HOT LIQUID GAS Expansion Condensation HOT LIQUID

The compressor pushes the liquid through the condenser coil at a high pressure (~ 10 atm). The liquid sprays through a throttling valve into the evaporation coil which is maintained by the compressor at a low pressure (~ 2 atm).



## Adiabatic demagnetization



Sample in flask B

- 1. Thermal contact (B filled), Field on.
  - → Xtal magnetized & cold
- 2. Thermal isolated (B pumped)
- 3. Field off
  - → Xtal demagnetized adiabatically
  - → Xtal cools down (entropy increases)

William F. Giauque (Nobel prize Chem. 1949) "for his contributions in the field of chemical thermodynamics, particularly concerning the behaviour of substances at extremely low temperatures"



## Adiabatic demagnetization



is removed

During adiabatic (better: isentropic) demagnetization the total entropy of the specimen is constant.

The entropy can flow into the spin system only from the system of lattice vibrations. The initial entropy of the lattice should be small in comparison with the entropy of the spin system in order to obtain significant cooling of the lattice – indeed,  $S_{lattice} \ll S_{spin}$  holds in the mK range.

The heat capacity of the spins in the two-state paramagnetic is *large* at low *T*: 1 cm<sup>3</sup> of iron ammonium salt at *T*=50 mK and *B* ~0.05 T has a heat capacity equal to 16 tons of lead (!) at the same *T*.

#### Temperatures attained:

- ~ 1 mK with electronic paramagnetic systems and
- ~ 1  $\mu$ K with nuclear paramagnetic systems.

## Intermetallic alloys



C. Poole, Superconductivity, Acad.Press 2007



Correlation sommerfeld coefficient & T<sub>c</sub> → Again density of states Fermi level

## Intermetallic compounds: The A15's

- A<sub>3</sub>B with
  A = Ti, V, Cr, Zr, Nd, Mo, Ta, W (transition metals)
  B = Sn, Ge, Ga, Si, ...
- Large number of SC's known (>60)
- Technologically relevant (superconducting magnets)

Tabelle 2.2	Supraleitende	Verbindungen	mit $\beta$ -Wolframstruktur	[5, 30].	• ИО
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Stoff	T <sub>c</sub> in K	$\lambda_L$ in nm	$\xi_{GL}$ in nm	B <sub>c2</sub> in T
V <sub>3</sub> Ge	6,0	65		
$V_3Ga^*$	14,2-14,6	65	4	23
V <sub>3</sub> Si	17,1	70	4	23
Nb <sub>3</sub> Sn	18,0	80	4	24
Nb <sub>3</sub> Ge	23,2	80	3	38 Duche

## Intermetallic compounds: The A15's



- Cubic crystal structure
- Like for alloys: Max Tc for N ~ 4.5 & 6.5 e/a (Matthias rule)
- Again correlation with Sommerfeld parameter Magn. Susceptibility

→ D(E<sub>F</sub>)

# Intermetallics: MgB<sub>2</sub>

#### Nagamatsu, Akimitsu Nature 2001

- Highest T<sub>c</sub> of conventional (e-ph) superconductors. 40 K
- Type II (H<sub>c1</sub>=2T, H<sub>c2</sub>=15-20T)
- Hexagonal crystal structure
- Material commercially available since 50's
- 2-band SC (2 bands cross E<sub>F</sub>)



## Intermetallics: MgB<sub>2</sub>

- σ- and π-bands crossing E<sub>F</sub>
   (Boron layer, analogy graphite)
- Two gaps!!

 $2\Delta_{\sigma} = 7 \text{ meV}$  $2\Delta_{\pi} = 2-4 \text{ meV}$ 

• 2-band BCS-like theory





## Fullerides



E. A. Rohlfing,<sup>a)</sup> D. M. Cox, and A. Kaldor J. Chem. Phys. 81 (7), 1 October 1984

### Nobelprize (Chem.) Fullerenes 1996



Fig. 1. Electronic absorption spectrum of: (a)  $C_{60}^+$  obtained by growing a neon matrix with a mass-selected ion beam; (b)  $C_{60}^+$  and  $C_{60}^-$  generated in a neon matrix upon NeI irradiation during deposition of  $C_{60}$  with neon.



Sheet: M. Grüninger

#### Derivation [edit]

An informal derivation follows. Suppose that f(x, y, z) = 0. Write z as a function of x and y. Thus the total differential dz is

$$dz = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x dy$$

Suppose that we move along a curve with dz = 0, where the curve is parameterized by x. Thus y can be written in terms of x, so on this curve

$$dy = \left(\frac{\partial y}{\partial x}\right)_z dx$$

Therefore the equation for dz = 0 becomes

 $0 = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x \left(\frac{\partial y}{\partial x}\right)_z dx$ 

Since this must be true for all dx, rearranging terms gives

$$\left(\frac{\partial z}{\partial x}\right)_y = -\left(\frac{\partial z}{\partial y}\right)_x \left(\frac{\partial y}{\partial x}\right)_z$$

Dividing by the derivatives on the right hand side gives the triple product rule

$$\left(\frac{\partial x}{\partial y}\right)_z \left(\frac{\partial y}{\partial z}\right)_x \left(\frac{\partial z}{\partial x}\right)_y = -1$$

Note that this proof makes many implicit assumptions regarding the existence of partial derivatives, the existence of the exact differential dz, the ability to construct a curve in some neighborhood with dz = 0, and the nonzero value of partial derivatives and their reciprocals. A formal proof based on mathematical analysis would eliminate these potential ambiguities.

#### Wikipedia